

NLSE Soliton Spectral Efficiency for Gaussian Input

Pavlos Kazakopoulos⁽¹⁾ and Aris L. Moustakas⁽²⁾

⁽¹⁾CS Research Foundation, Amsterdam

⁽²⁾Physics Department, National and Kapodistrian University of Athens

Abstract

- Imminent “**Capacity Crunch**” necessitates new look at channel nonlinearities in Fiber Optic Communications.
- The channel nonlinearities can be adequately described by the **non-linear Schrodinger equation (NLSE)**, which is integrable.
- **In this work** we take advantage of exact results when the input signal distribution is Gaussian to derive an **analytic lower bound for the spectral efficiency** of solitonic transmission in the low noise limit, using scattering data parameters which can be evaluated numerically.

Channel Model

- Single Fiber Channel NLSE:

$$i \frac{\partial u}{\partial z} + \frac{\partial^2 u}{\partial t^2} + 2|u|^2 u = \varepsilon \sum_{k=1}^M f_k(t) \delta(z - kL)$$

- Normalized units:

$$z \rightarrow \frac{z}{L_s} \quad t \rightarrow \frac{t}{t_s} \quad u \rightarrow \frac{u}{\sqrt{P t_s}}$$

$$t_s = \sqrt{\frac{\beta_2}{\gamma P}} \quad L_s = \frac{2}{P t_s} \quad \varepsilon^2 = \frac{\sigma^2}{P t_s} = \frac{\alpha h \nu K}{P t_s}$$

- Bandwidth: $(-B, B)$ $L_d^{-1} = \beta_2 B^2$

- Noise Model: M amplifiers, L apart, with white noise:

$$E[f_k(t) \bar{f}_k(t')] = \delta(t - t')$$

Inverse Scattering Transform

- Exploit integrability of NLSE:

- Given input signal $u(t)$ calculate scattering data of associated linear Zakharov-Shabat equation

$$\begin{bmatrix} i\partial_t & u \\ -\bar{u} & -i\partial_t \end{bmatrix} \begin{pmatrix} \Psi_1 \\ \Psi_2 \end{pmatrix} = \lambda \begin{pmatrix} \Psi_1 \\ \Psi_2 \end{pmatrix}$$

- 4 degrees of freedom per soliton

$$\eta_n, \xi_n, \log b_n = \eta_n X_n + i\phi_n$$

- Jost functions $\Psi_1(t)$, $\Psi_2(t)$ are an important byproduct of the calculation

- **Basic Benefit:** In the absence of noise, scattering data have no (or trivial) spatial propagation

Effect of Noise

- Noise shifts of the scattering data:

$$\delta\lambda_n = \varepsilon \frac{\int (f_k \Psi_{1n}^2 + \bar{f}_k \Psi_{2n}^2)}{\int \Psi_{1n} \Psi_{2n}}$$

$$\delta \log b_n = -i2\lambda_n L \sum_k \delta\lambda_{nk} (M - k)$$

$$+ i \frac{\varepsilon b_n}{a'(\lambda_n)} \sum_k \int (f_k \partial \Psi_{1n} \Psi_{1n} - \bar{f}_k \partial \Psi_{2n} \Psi_{2n})$$

$$\partial \Psi_{in} = \left[\partial_\lambda \Psi_i - b_n^{-1} \partial_\lambda \phi_i \right]_{\lambda=\lambda_n} \pm \log b_n \Psi_i$$

- In scattering data, first term corresponds to Gordon-Haus effect: Random shift in velocity
- Second term: “intrinsic” shift

Soliton Information Theory

- Input Distribution: From signal to scattering data

$$u(t) \rightarrow \{\eta, \xi, X, \phi\}$$

$$P(u) \rightarrow P(\eta, \xi, X, \phi)$$

- Gaussian Noise Model:

$$f(t) \rightarrow \{\delta\eta, \delta\xi, \delta X, \delta\phi\}$$

- For low noise, effectively AWGN channel (but with **colored noise**)

- From Sequences to Sets:

- Solitons are indistinguishable (Gordon-Haus effect mixes any sequence)
- Thus, for X iid:

$$H_{set}(X) = \sum_{n=1}^N H(x_n) - \log N!$$

- Noise Entropy: Use of Permanent

$$H_{set}(Y | X) = - \sum_p P(Y | X_p) \log \left(\sum_{p'} P(Y | X_{p'}) \right)$$

- Low noise limit: Only one permutation probable. Thus:

$$H_{set}(Y | X) \approx N \left(\log \sqrt{2\pi e} + \frac{1}{2N} \log \det \Sigma \right)$$

- Σ 4x4 block matrix of NxN correlations

Gaussian Input

- Assume Gaussian input distribution

$$P(\{u(t)\}) \propto e^{-\frac{1}{2D} \int |u(t)|^2 dt}$$

- Correlations:

$$E[u(t) \bar{u}(t')] = 2D \delta_b(t - t')$$

- From dimensional analysis: $D = \frac{1}{4Bt_s}$

- Why?

- Exact distribution:

$$P(\eta, \xi) = \frac{2}{D} \frac{\coth\left(\frac{\eta}{D}\right) - 1}{\sinh^2\left(\frac{\eta}{D}\right)}$$

- Conserved during propagation
- $\Psi_1(t)$, $\Psi_2(t)$ localized (positive Lyapunov exponent), thus scattering data (nearly) independent and (a.s.) solitonic.
- Uniform distribution in X and ϕ
- D only remaining scale in ZS system: Hence $t \rightarrow tD^{-1}$ $u \rightarrow Du$ $\lambda \rightarrow \lambda D$

- Number of solitons

$$E \left[\int |u|^2 dt \right] = 4DBT = 4NE[\eta] = 4ND$$

- Hence $N = BT$

Spectral Efficiency

- 2BT complex degrees of freedom
- N = BT solitons with 4 real degrees of freedom

$$SE \geq \frac{N}{2BT} \left(\begin{array}{l} H(\eta) + H(\xi) + H(X) + H(\phi) \\ - \log \left(\frac{N}{e} \right) \\ - \log \sqrt{2\pi e} - \frac{1}{2N} \log \det \Sigma \end{array} \right)$$

- After calculating input distributions and reintroducing units

$$SE \geq \log \left(\frac{Pe}{B\sigma^2 M} \right) - \frac{1}{4N} E[\log \det S]$$

- Symmetric S matrix:

$$S = \begin{bmatrix} R_{\eta\eta} & R_{\eta\xi} & F_1 + q^2 F_2 & D_1 + q^2 D_2 \\ \dots & R_{\xi\xi} & G_1 + q^2 G_2 & A_1 + q^2 A_2 \\ \dots & \dots & C_1 + q^2 C_2 + q^4 C_3 & B_1 + q^2 B_2 + q^4 B_3 \\ \dots & \dots & \dots & P_1 + q^2 P_2 + q^4 P_3 \end{bmatrix}$$

- Submatrices are O(1) correlation matrices

$$q = \frac{1}{4} \sqrt{\frac{L_{tot}}{\beta_2}} \frac{\gamma P}{B} = \frac{\sqrt{L_{tot} L_d}}{2L_{nl}}$$

- When $q \ll 1$ nonlinearity weak and Shannon limit correct

$$SE \approx \log(SNR)$$

- When $q \gg 1$ nonlinearity is strong (hence Gordon-Haus effect dominates)

$$SE \approx \log \left(\frac{\beta_2 B}{ML_{tot} \sigma^2 \gamma^2 P} \right) + O(1)$$

Conclusions

- Analytic expression of soliton spectral efficiency in terms of fundamental parameters of fiber.
- Expressed in terms of correlations functions, written as integrals over Jost functions that can be averaged over realizations.
- Expected to self-average in large system-size limit.
- Crossover between linear AWGN and non-linear regime determined by scalar parameter q.
- Methodology can also be applied to defocusing NLSE and multimode integrable systems.

